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Measurement of the $T+T$ neutron spectrum using the National Ignition Facility

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Neutron time-of-flight spectra from inertial confinement fusion experiments with tritium-filled targets have been measured at the National Ignition Facility. These spectra represent a significant improvement in energy resolution and statistics over previous measurements, and afford the first definitive observation of a peak resulting from sequential decay through the ground state of ${}^5\text{He}$ at low reaction energies $E_{\text{c.m.}} \lesssim 100$ keV. To describe the spectrum we have developed an R -matrix model that accounts for interferences from fermion symmetry and intermediate states, and show these effects to be non-negligible. We also find the spectrum can be described by sequential decay through $\ell = 1$ states in ${}^5\text{He}$, which differs from previous interpretations.

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Investigations of the neutron spectrum from the low energy $T(t, 2n)\alpha$ reaction [1] have spanned six decades, driven by interests ranging from fusion research [2–4] to the neutron-neutron scattering length [5]. The spectrum provides information about interactions among final-state particles, while the $T(t, 2n)\alpha$ yield extracted from spectral measurements assists efforts to understand thermonuclear environments created in the laboratory. Reaction yields test the modeling of inertial confinement fusion, for example, in experiments designed to measure mixing between the ablator and fuel [6, 7]. Also, since the particle spectrum results from three-body decay its sensitivity to different breakup mechanisms has long been of interest [8]. Energy resolution, backgrounds, and statistics have presented major experimental challenges in the past, and a modern shortage of tritium beams has required developing alternative approaches to study the spectrum. Still, large uncertainties remain in the spectrum, and disparate accounts of its underlying breakup processes exist [3–5, 8].

The particle spectrum from the $T(t, 2n)\alpha$ reaction has been described previously by statistical and/or sequential processes [3–5, 8–12]. Statistical decays uniformly populate the available phase space, producing a neutron spectrum in the shape of an ellipse. Sequential decays proceed through intermediate states formed by $n\text{-}\alpha$ (${}^5\text{He}$) or $n\text{-}n$ (correlated di-neutron) interactions. (Note: the distinction between statistical and sequential mechanisms can become ambiguous, since broad intermediate states can also fill phase space uniformly.) When transitions

are made through ${}^5\text{He}$, the neutrons emitted first are spread over the decay width of the intermediate state by energy-momentum conservation. The distribution of neutrons from ${}^5\text{He}$ is further broadened by recoil, with a shape reflecting the angular correlation between neutrons. For correlated di-neutron decays, the phase-space distribution of energies is modified by an s -wave singlet interaction between final-state neutrons.

Discussions in the literature of sequential decay through ${}^5\text{He}$ must include the interference effects arising from a final state with identical fermions. These order-of-emission effects have been well-established with studies involving the 3α breakup of ${}^{12}\text{C}$ [13–15] and also have been applied to the sequential proton decay of ${}^6\text{Be}$ [16]. The coherence between accessible states in ${}^5\text{He}$ should be considered as well. For the ground $J^\pi = 3/2^-$ and first-excited $1/2^-$ states in ${}^5\text{He}$ this produces an interference in the angular correlation between neutrons and, by extension, the spectrum. We [17] have formulated sequential decay processes within the R -matrix framework [18, 19], where both interferences can be incorporated. This model reveals substantial changes occurring from antisymmetrization.

Experimental studies of the neutron spectrum have taken a variety of approaches to investigate the $T(t, 2n)\alpha$ reaction, covering a range in center-of-mass energies from $E_{\text{c.m.}} = 750$ keV [9] down to nearly zero energy [2, 3]. In addition to earlier accelerator work [5, 8, 9, 12], fusion experiments based on muon catalyzed [2, 3] and inertial confinement [4] have been conducted. Although the pre-

cise details of detectors differed among experiments, most have deduced the neutron spectrum from either proton or deuteron recoils [2–4, 8, 9]. It is interesting to note that only three groups have reported their experimental energy resolution (full-width at half-maximum): 0.5 [12], 1.3 [4], and 2.8 MeV [9]. The work presented in this Letter has combined the inertial confinement approach with time-of-flight technique to perform measurements resolved to 280 keV [20] – a factor of ≈ 2 improvement over any previous experiment.

The first reported measurement of the neutron spectrum ($E_{\text{c.m.}} = 110$ keV) [8] was interpreted as a statistical distribution, accompanied by a much smaller ($\approx 10\%$) ${}^5\text{He}$ ground-state contribution. Wong *et al.* [12] later determined at slightly higher energies (250 keV) that 70% of decays proceeded through correlated di-neutron emission, with the remainder split between the ground (20%) and first-excited (10%) states of ${}^5\text{He}$. As stipulated in their paper the branching ratios were directly tied to an assumption of a strong di-neutron channel. The predominance of that channel is contrary to the findings of Poutissou *et al.* [5] (20 keV), who stated any n - n correlations were masked by those from n - α , and Matsuzaki *et al.* [3] (≈ 0 keV), who emphasized the importance of n - α interactions. Yet the conclusions of both groups were strongly model dependent as no definitive signature of n -

α interactions was visible in their data, such as the peak at 8.7 MeV in the spectra of Refs. [8, 12] from sequential decays through the ground state of ${}^5\text{He}$. The significance of these interactions at low reaction energies was further brought into question by a recent Letter [4] that reported no observable n - α effects at $E_{\text{c.m.}} = 23$ keV. Casey *et al.* [4] pointed out the diminishing strength of the ${}^5\text{He}$ ground-state peak with reaction energy, from Wong *et al.* to Allen *et al.* [8] to Poutissou *et al.*, who also reported no appreciable contribution from this channel. In addition, the correlated di-neutron distribution of Refs. [11, 12] was found to describe the measurements of Casey *et al.* Utilizing a new experimental capability, we report in this Letter the first direct observation of the peak from the ${}^5\text{He}$ ground-state channel at low energies. Further, with the R -matrix model, we find the main features in these new measurements can be described by sequential decay through $\ell = 1$ states in ${}^5\text{He}$.

The energy distribution of particles emitted by reactions in unbound states can be described with R -matrix methods, as given by Barker [19]. But when a reaction leads to a final state containing identical particles, the standard formalism must be adapted to treat an amplitude with exchange symmetry. For the low energy $T(t, 2n)\alpha$ reaction, coupling tritons in an s wave to form a 0^+ state, the amplitude can be generalized as

$$\mathcal{M}_{\nu_1\nu_2}^{\ell J} = \sum_{m_\ell m'_\ell m} \frac{(-1)^{J+m}}{\sqrt{2J+1}} (\ell m_\ell \frac{1}{2} \nu_1 | J m) (\ell m'_\ell \frac{1}{2} \nu_2 | J - m) Y_{\ell m_\ell}(\hat{p}_1) Y_{\ell m'_\ell}(\hat{p}_{23}) \frac{e^{-i(\Phi_1 + \Phi_{23})} \sqrt{\frac{P_1 P_{23}}{p_1 p_{23}}} \sum_\lambda \frac{A_{J\lambda} \gamma_{J\lambda}}{E_{J\lambda} - E_{23}}}{1 - (B_J - S_{23} + i P_{23}) R_{J\lambda}}. \quad (1)$$

The numerical subscripts in Eq. (1) reflect the decay process: emission of neutron 1 and the subsequent disintegration of ${}^5\text{He}$ to neutron 2 and α particle 3. For a given J^π in ${}^5\text{He}$ and orientations ν of emitted neutrons, summations are performed over possible angular momentum quantum numbers. Clebsch-Gordan coefficients along with spherical harmonics $Y_{\ell m_\ell}$, which depend on projections of momenta p in the three-body \hat{p}_1 and recoil center-of-mass \hat{p}_{23} systems, describe the angular correlation between neutrons. Only the ground $3/2^-$ and first-excited $1/2^-$ states of ${}^5\text{He}$ are directly accessible from reaction energies relevant here, and for both cases $\ell = 1$ neutrons are ejected in each stage of the breakup. Energy levels $E_{J\lambda}$, reduced-width amplitudes $\gamma_{J\lambda}$, boundary values B_J , and feeding factors $A_{J\lambda}$ are the R -matrix parameters contained above, and all except the $A_{J\lambda}$ factors can be constrained with high precision n - α scattering data [21]. The remaining terms are calculable for a given channel radius and partial wave: the hard sphere phase shift Φ and the penetration P and shift S factors.

Exchanging neutron indices in the Eq. (1) yields a second expression, which is necessary to construct the antisymmetric amplitude. The energy distribution is obtained by summing the amplitude over J , squaring the magnitude, then summing over ν_1 and ν_2 , and multiplying by the three-body phase-space factor [22]. Interferences from symmetry and intermediate states originate from cross terms in the squared amplitude. Lastly, to describe kinematically incomplete measurements, such as the spectra considered here, an integration is performed over angles of the unobserved neutron. Fig. 1 presents spectra that result from the procedure outlined above. The spectrum calculated for the $1/2^-$ partial wave (middle panel) illustrates the importance of antisymmetry, which leads to a $\approx 30\%$ contribution.

Inertial confinement experiments were conducted at the National Ignition Facility. Here, temporally-shaped laser beams were used to generate a soft x-ray field that ablatively imploded a spherical plastic capsule filled with 99.9% tritium gas (0.1% deuterium). Using time-of-flight detectors, an ion temperature of $k_B T = 3.3(3)$ keV [or

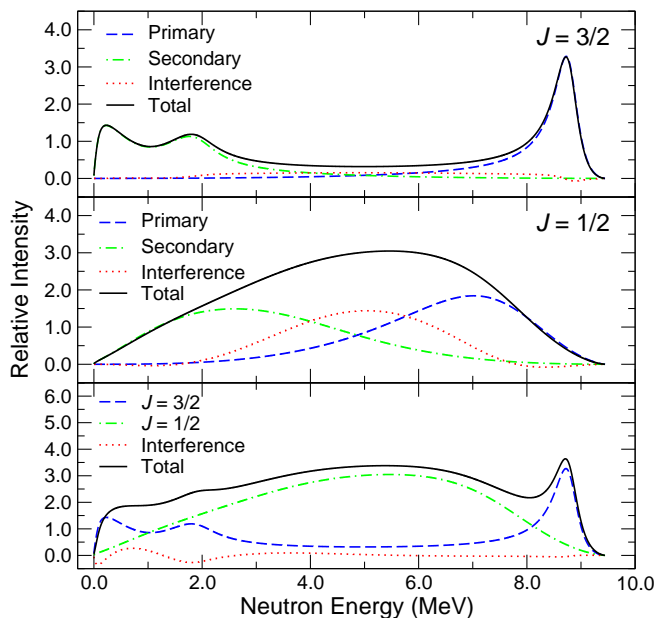


FIG. 1. (color online) R -matrix spectra resulting from the fitted feeding factors are shown above. Decays through partial waves of the ground and first-excited states of ^5He are shown in the top and middle panels, respectively. Each partial wave contains an interference term from antisymmetrization, which is proportional to the overlap between primary and secondary distributions. For the $J = 1/2^-$ partial wave, the interference term makes a $\approx 30\%$ contribution. The bottom panel displays the coherent sum of both partial waves. An interference created within the neutron angular correlation manifests itself in the spectrum. The feeding factors determine the sign of this interference.

an effective energy for the $T(t, 2n)\alpha$ reaction of $E_{\text{c.m.}} = 16(1)$ keV] within the burn volume was extracted from thermal broadening of nearly monoenergetic 14-MeV-neutrons produced by the $T(d, n)\alpha$ reaction. A total neutron yield in excess of 10^{13} was produced over the ≈ 200 -ps duration of thermonuclear reactions, affording excellent statistics and timing. Measurements have also been performed with a capsule containing a deuterium-to-tritium ratio of 3:1. These data serve to establish the neutron-induced background from residual deuterium in the capsules, and show a relatively flat baseline several decades below the preceding $T(d, n)\alpha$ peak (see Fig. 2, bottom panel).

Neutron spectra were recorded by two 10-cm-by-5-cm xylene scintillators, positioned over 20 m from the target chamber center along separate lines of sight [23]. Each line was doubly collimated with 1 m of steel: the first collimator defined a 8.3-cm-diameter opening on the detector face and the second was designed to clean up after the first. A 1.3-cm thick attenuator (95% W, 3% Ni and 2% Fe) covered each line of sight to suppress the photon background created during implosions. Attenuation and scattering of the neutron spectrum from interven-

ing material has been taken into account with MCNP [24] simulations, and corrections were $\lesssim 10\%$ above 2 MeV.

Implosions are singular events which generate tremendous neutron fluxes, and so the standard methods of pulse counting are not applicable. Detectors were operated in current mode and their signals were digitized with 1-GHz oscilloscopes. Interpreting current mode data requires knowledge of the impulse response functions (IRF) from detectors. As the light decay-time of the scintillator is rather insensitive to particle type, short pulses (100 ps) of x rays have been used to measure the IRFs. Further refinements have been made to account for neutron transport in the scintillator. Constructed IRFs have been validated over many measurements of the $T(d, n)\alpha$ peak, in which good fits to data ($\chi^2_\nu \approx 1$) were obtained with the expected distribution from thermal broadening [25]. Careful attention has also been given to describing light production in the scintillator over observed neutron energies. Consistency checks of the scintillator response have been performed *in situ* with activation foils, which agree with the scintillators to 8%.

The R -matrix spectrum was fitted simultaneously to data from both detectors. This method combines neutron statistics and averages over systematic differences between measurements. The fitting algorithm first thermally broadens a trial spectrum then adjusts for attenuation/scattering and scintillation response. Next, for each detector, the algorithm transforms the modified spectrum from energy into time of flight and convolves with the IRFs. The top panel of Fig. 2 illustrates these experimental corrections. A χ^2 minimization of the resulting spectrum to data determines free parameters with the MINUIT2 package [26]. Four feeding factors $A_{J\lambda}$ were allowed to vary in the fit; all remaining parameters were fixed to values reported by Stambach and Walter [21].

Data corresponding to neutron energies $\lesssim 1$ MeV were excluded from the analysis presented in Fig. 2 because of limited detector sensitivity. Digitizer noise from oscilloscopes represented the main source of statistical uncertainty in the data. Fluctuations due to neutron events were much smaller, as $\approx 10^6$ neutrons interacted with the scintillator to produce the signal in the region of interest. Systematic uncertainties have been investigated for ion temperature, attenuation/scattering, region of interest, and scintillator response. These uncertainties become more important for data below 4.5 MeV.

Separate analyses were undertaken for each systematic error at its 1σ deviation, and the sum of these individual errors (16%) was added in quadrature with uncertainties determined from the fitting routine (1%) to give a total error of 16%. The best fit results in a $\chi^2 = 1440$ for 660 data points and 4 free parameters ($\chi^2_\nu = 2.2$), and comparable values of χ^2_ν result for each detector. This fit determines branching ratios between the $3/2^-$ and $1/2^-$ states of 29 ± 16 and $71 \mp 16\%$, respectively. The error

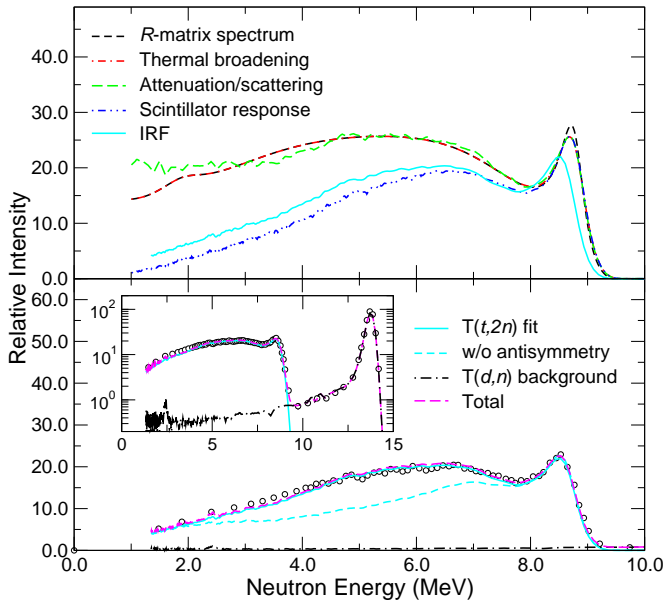


FIG. 2. (color online) Experimental corrections applied to the modeled spectrum, as discussed in the text, are displayed in successive order (top panel). The combined fit is plotted over energy-binned data from the 22-m detector (bottom panel), in the region of interest and over the full range of neutron energies. Statistical errors are approximately the size of the data points. Transitions through the ground state of ${}^5\text{He}$ produce a peak at ≈ 8.7 MeV; the broad distribution peaked around 6 MeV can largely be described by the first-excited state in ${}^5\text{He}$. For reference the short-dashed (cyan) curve shows the effect of not including antisymmetry. The background from deuterium contamination, which is illustrated better in the inset, has a minimal effect in the region of interest.

in the ratio is anti-correlated to maintain unity. Note that the ratio is largely insensitive to the interference between partial waves, as can be seen in the bottom panel of Fig. 1. While differences between the fitted curve and data appear in several locations, the model describes the main features in the spectrum quite well, particularly in the region of the spectrum that is least susceptible to systematic errors ≥ 4.5 MeV.

New spectra from inertial confinement fusion experiments with enriched tritium targets have been measured at the National Ignition Facility. These measurements provide conclusive evidence for $n\text{-}\alpha$ interactions in the $T + T$ neutron spectrum produced by thermonuclear environments, through the first direct observation of a neutron peak from transitions via the ground state of ${}^5\text{He}$. A description of the data was obtained with an R -matrix model that accounts for the coherence between intermediate states and fermion symmetry, which was shown in this Letter to be non-negligible. Fits of the model to the spectrum show a large contribution from the $1/2^-$ partial wave in ${}^5\text{He}$, which differs from the dominant role found for either correlated di-neutron emission [4] or statistical

decay [8]. Future work will investigate the possible role of di-neutron and/or sequential s -wave emission within the framework of the new R -matrix approach.

Although a reasonable description of the data was obtained with the R -matrix model, we acknowledge the present analysis has limitations as it represents a two-body approximation to what is strictly a three-body problem. Also, the interpretation of the neutron singles spectrum can be subject to ambiguities that only kinematically complete measurements distinguish. We hope these new data encourage further work in theory and experiment, and improve evaluations of fusion experiments.

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- [1] We consider the low energy $T(t,2n)\alpha$ reaction over a region $E_{c.m.} \lesssim 100$ keV where the s -wave channel is dominant and the astrophysical S -factor is flat [27, 28].
- [2] W. H. Breunlich *et al.*, Muon Catal. Fus. **1**, 121 (1987).
- [3] T. Matsuzaki *et al.*, Phys. Lett. B **557**, 176 (2003).
- [4] D. T. Casey *et al.*, Phys. Rev. Lett. **109**, 025003 (2012).
- [5] R. Larose-Poutissou and H. Jeremie, Nucl. Phys. **A218**, 559 (1974).
- [6] D. C. Wilson *et al.*, Phys. Plasmas **18**, 112707 (2011).
- [7] V. A. Smalyuk, (manuscript in preparation).
- [8] K. W. Allen *et al.*, Phys. Rev. **82**, 262 (1951).
- [9] S. J. Bame and W. T. Leland, Phys. Rev. **106**, 1257 (1957).
- [10] N. Jarmie and R. C. Allen, Phys. Rev. **111**, 1121 (1958).
- [11] B. Lacina, J. Ingley, and D. W. Dorn, UCRL Report No. 7769 (Lawrence Radiation Lab., Univ. of California, Livermore, 1965).
- [12] C. Wong, J. D. Anderson, and J. W. McClure, Nucl. Phys. **71**, 106 (1965).
- [13] P. Swan, Rev. Mod. Phys. **37**, 336 (1965).
- [14] D. P. Balamuth, R. W. Zurmühle, and S. L. Tabor, Phys. Rev. C **10**, 975 (1974).
- [15] H. O. U. Fynbo *et al.*, Phys. Rev. Lett. **91**, 082502 (2003).
- [16] D. F. Geesaman *et al.*, Phys. Rev. C **15**, 1835 (1977).
- [17] C. R. Brune, (manuscript in preparation).
- [18] A. Lane and R. Thomas, Rev. Mod. Phys. **30**, 257 (1958).
- [19] F. C. Barker, Aust. J. Phys. **41**, 743 (1988), we replace $G_{\lambda bc}^{1/2}$ by $P^{1/2}A_\lambda$ in his Eq. (27).
- [20] The intrinsic detector resolution for 8.7-MeV neutrons (180 keV) is thermally broadening by 220 keV, resulting in an effective resolution of 280 keV.
- [21] T. Stammach and R. L. Walter, Nucl. Phys. **A180**, 225 (1972), we fix our R -matrix parameters for $E_{J\lambda}$, $\gamma_{J\lambda}$, B_J and channel radius to the values listed in their Table 2.
- [22] G. G. Ohlsen, Nucl. Instr. and Meth. **37**, 240 (1965), the three-body phase space term is given by his Eq. (40).
- [23] J. A. Caggiano, (manuscript in preparation).
- [24] R. Forster *et al.*, Nucl. Instrum. Meth. **B213**, 82 (2004).

- [25] L. Ballabio, J. Kallne, and G. Gorini, Nucl. Fusion **38**, 1723 (1998).
- [26] F. James and M. Winkler, *Minuit2: Function Minimization and Error Analysis*, CERN, Geneva (2004).
- [27] R. E. Jarmie and R. E. Brown, Nucl. Instr. and Meth. **B10/11**, 405 (1985).
- [28] C. Angulo and P. Descouvemont, Nucl. Phys. **A639**, 733 (1998).